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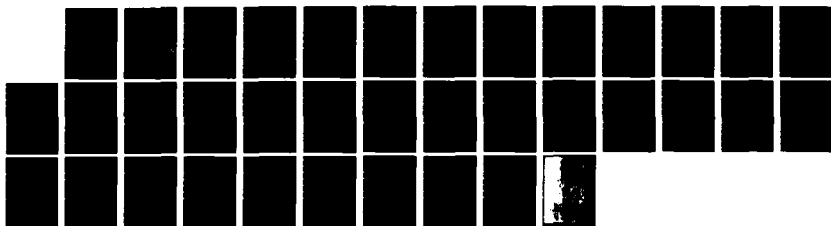
SIGNIFICANCE POINTS FOR SOME TESTS OF UNIFORMITY ON THE
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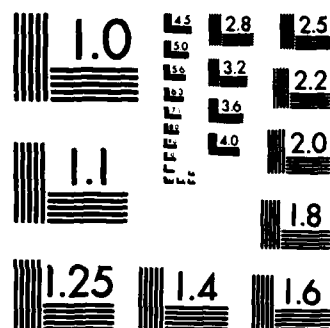
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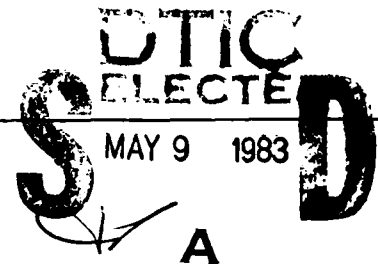
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ITEM #2c, CONTINUED:

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(2) presenting ^{SOME} ~~some~~ percentage points for selected small and moderate sample sizes obtained by Monte-Carlo methods;

(3) evaluating numerically the cumulative distribution functions and significance points of the limiting distributions via the Laguerre transform method (Keilson and Nunn (1979), Keilson, Nunn and Sumita (1981), and Sumita (1981)).

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OF UNIFORMITY ON THE SPHERE

by

J. Keilson
D. Petronidas
U. Sumita
J. Wellner

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Significance points for some tests of uniformity
on the sphere

J. Keilson, D. Petronidas, U. Sumita, J. Wellner

Abstract

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- (1) finding the exact distributions of the statistics of Beran and Giné for $n = 2$;
- (2) presenting some percentage points for selected small and moderate sample sizes obtained by Monte-Carlo methods;
- (3) evaluating numerically the cumulative distribution functions and significance points of the limiting distributions via the Laguerre transform method, (Keilson and Nunn (1979), Keilson, Nunn and Sumita (1981), and Sumita (1981)).

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§1. Introduction

Beran (1968) and Giné (1975) have proposed several omnibus tests for uniformity on the unit sphere $S = \{\underline{x} \in R^3: |\underline{x}| = 1\}$ in three-dimensional Euclidean space. These tests are consistent against all alternatives and are locally most powerful for specific alternatives. Beran and Giné have shown that the limiting distributions of these statistics, under the null hypothesis of uniformity, are those of weighted sums of independent Chi-square variables. While Prentice (1975) has applied the methods of Zolotarev (1961) and Hoeffding (1964) to provide approximate percentage points for the limiting distributions, no tables of the limiting distributions, percentage points thereof, or finite sample distributions or percentage points have been available.

Our purpose here is to fill this gap by:

- (1) finding the exact distributions of the statistics of Beran and Giné for $n = 2$;
- (2) presenting some percentage points for selected small and moderate sample sizes obtained by Monte-Carlo methods;
- (3) using the Laguerre transform method - Keilson and Nunn (1979), Keilson, Nunn and Sumita (1981), and Sumita (1981) to compute the cumulative distribution functions and significance points of the limiting distributions.

In Section 2, we summarize the result on the limiting distributions due to Beran and Giné and give the exact distributions for $n = 2$. Section 3 contains the Monte-Carlo results for finite sample sizes and description of the methods used. We discuss, in Section 4, the numerical procedure for

evaluating the limiting distributions via the Laguerre transform method. An application is given in Section 5 where we test uniformity of orientation of dendritic fields in the retinas of cats subject to controlled visual environments. The numerical results are summarized in Section 6 in tables and graphs.

§2. The statistics: limiting distributions and exact distributions for $n = 2$

Let $\underline{X}_1, \dots, \underline{X}_n$ be independent and identically distributed unit vectors in R^3 with distribution ν (so $\nu(S) = 1$, where $S = \{\underline{x} \in R^3: |\underline{x}| = 1\}$). Let $\theta_{ij} \equiv \hat{\underline{X}}_i \cdot \hat{\underline{X}}_j \equiv \arccos(\underline{X}_i \cdot \underline{X}_j)$ = the angle between \underline{X}_i and \underline{X}_j for $i, j = 1, 2, \dots, n$.

For testing the null hypothesis that ν is the uniform distribution on S , Beran (1968) and Giné (1975) have suggested the statistics

$$Y_n^{\text{odd}} \equiv n^{-1} \sum_{i=1}^n \sum_{j=1}^n \{1 - (2/\pi)\theta_{ij}\} \quad (2.1)$$

$$Y_n^{\text{even}} \equiv n^{-1} \sum_{i=1}^n \sum_{j=1}^n \{\frac{1}{2} - (2/\pi)\sin \theta_{ij}\} \quad (2.2)$$

$$Y_n \equiv Y_n^{\text{odd}} + Y_n^{\text{even}} \quad (2.3)$$

Giné proposes Y_n^{even} and finds that Y_n^{odd} is simply Beran's form of Aine's statistic: $Y_n^{\text{odd}} = 4T_n$ (Beran, 1968). It is also known that $Y_n^{\text{odd}} = 4A_{2,n}$ (Prentice, 1978) and $Y_n^{\text{even}} = G_{2,n}$ (Prentice).

Let $\{Z_j\}_{j=1}^{\infty}$ be independent Chi-square random variables with $2j+1$ degrees of freedom (so $Z_{2k-1} \sim \chi_{4k-1}^2$, $Z_{2k} \sim \chi_{4k+1}^2$). Set

$$a_{2k-1}^2 = (2k-1)^{-2} [(\frac{1}{2})_k / k!]^2 \quad (2.4)$$

$$a_{2k}^2 = (2k-1)^{-1} (2k+2)^{-1} [(\frac{1}{2})_k / k!]^2 \quad (2.5)$$

where $(\frac{1}{2})_k = \frac{1}{2} \cdot (\frac{1}{2} + 1) \cdot \dots \cdot (\frac{1}{2} + k-1)$ and define

$$Y^{\text{odd}} = \sum_{k=1}^{\infty} a_{2k-1}^2 Z_{2k-1} \quad (2.6)$$

$$Y^{\text{even}} = \sum_{k=1}^{\infty} a_{2k}^2 Z_{2k} \quad (2.7)$$

$$Y = Y^{\text{odd}} + Y^{\text{even}} \quad (2.8)$$

Theorem 1 (Giné): If ν is the uniform distribution on S , then

$$\lim_{n \rightarrow \infty} P[Y_n^{\text{odd}} \leq y] = P[Y^{\text{odd}} \leq y] ,$$

$$\lim_{n \rightarrow \infty} P[Y_n^{\text{even}} \leq y] = P[Y^{\text{even}} \leq y] ,$$

and $\lim_{n \rightarrow \infty} P[Y_n \leq y] = P[Y \leq y]$ for all $y \in \mathbb{R}$.

Prentice (1978) uses the methods of Zolotarev (1961) and Hoeffding (1964) to provide approximate percentage points for these limiting distributions. We will compute the distributions with precision via the Laguerre transform method [Keilson and Nunn (1979), Keilson, Nunn and Sumita (1981), and Sumita (1981)] in Section 4 and present tables and graphs in Section 6.

We now consider the distributions of Y_n^{odd} , Y_n^{even} and Y_n for $n = 2$.

Theorem 2: If ν is the uniform distribution on S , then

$$P[Y_2^{\text{odd}} \leq y] = \frac{1}{2}\{1 + \cos[\frac{\pi}{2}(y-2)]\} , \quad 0 \leq y \leq 2 ,$$

$$P[Y_2^{\text{even}} \leq y] = \cos\{\arcsin[\frac{\pi}{2}(1-y)]\} , \quad 1 - \frac{2}{\pi} \leq y \leq 1 ,$$

$$P[Y_2 \leq y] = \frac{1}{2}\{1 + \cos[g^{-1}(\frac{\pi}{2}(3-y))]\} , \quad 0 \leq y \leq 3 ,$$

where $g(t) \equiv t + \sin t$, $0 \leq t \leq \pi$, and g^{-1} denotes the inverse of g .

Proof: First note that under uniformity $P[\theta_{12} \leq y] = (4\pi)^{-1} \int_0^{2\pi} \int_0^y \sin \theta_{12} d\theta_{12} d\theta_{21} = \frac{1}{2}(1 - \cos y)$, $0 \leq y \leq \pi$. Then, writing $Y_2^{\text{odd}} = \frac{1}{2}\{1 + 1 + (1 - (2/\pi)\theta_{12}) + (1 - \frac{2}{\pi}\theta_{21})\} = 2 - (\frac{2}{\pi})\theta_{12}$, $Y_2^{\text{even}} = \frac{1}{2}\{\frac{1}{2} + \frac{1}{2} + (\frac{1}{2} - \frac{2}{\pi}\sin \theta_{12}) + (\frac{1}{2} - \frac{2}{\pi}\sin \theta_{21})\} = 1 - (\frac{2}{\pi})\sin \theta_{12}$, and $Y_2 = 3 - (\frac{2}{\pi})(\theta_{12} + \sin \theta_{12})$, the stated distributions are easily obtained by straightforward computation. \square

For future reference, we record some moments in the following Proposition:

Proposition 3. Under uniformity we have

$$E(Y_n^{\text{odd}}) = E(Y^{\text{odd}}) = 1$$

$$E(Y_n^{\text{even}}) = E(Y^{\text{even}}) = \frac{1}{2}$$

$$E(Y_n) = E(Y) = \frac{3}{2}$$

while

$$\text{Var}(Y_n^{\text{odd}}) = (1 - \frac{1}{n})(2 - \frac{16}{\pi^2}) \xrightarrow{n \rightarrow \infty} (2 - \frac{16}{\pi^2}) = .37886 \dots = \text{Var}[Y^{\text{odd}}] ,$$

$$\text{Var}(Y_n^{\text{even}}) = (1 - \frac{1}{n})(\frac{16}{3\pi^2} - \frac{1}{2}) \xrightarrow{n \rightarrow \infty} (\frac{16}{3\pi^2} - \frac{1}{2}) = .04038 \dots = \text{Var}[Y^{\text{even}}] ,$$

and

$$\text{Var}(Y_n) = (1 - \frac{1}{n})(\frac{3}{2} - \frac{32}{3\pi^2}) \xrightarrow{n \rightarrow \infty} (\frac{3}{2} - \frac{32}{3\pi^2}) = .41924 \dots = \text{Var}[Y] .$$

Furthermore, the third and fourth cumulants of Y^{odd} and Y^{even} are given by

$$K_3(Y^{\text{odd}}) = .375219 \dots , \quad K_4(Y^{\text{odd}}) = .56252 \dots ,$$

$$K_3(Y^{\text{even}}) = .0098016 \dots , \quad K_4(Y^{\text{even}}) = .00366374 \dots .$$

Proof:

The finite sample means and variances are easily obtained by elementary methods upon noting that θ_{ij} and $\theta_{ij'}$ are independent for $j \neq j'$; the asymptotic means and variances follow immediately by letting $n \rightarrow \infty$. The cumulants of Y^{odd} and Y^{even} are easily computed using the following easily

derived formula for the m^{th} cumulant K_m of $\sum_{j=1}^{\infty} \beta_j x_{f_j}^2$:

$$K_m = 2^{m-1} (m-1)! \sum_{j=1}^{\infty} f_j \beta_j^m . \quad \square$$

§3. Monte-Carlo simulations

In our simulations, a function subprogram called RAND (University of Rochester file #311.7.500, Computer Center) and the IBM/360 computer are used to generate uniformly distributed random numbers. The method of Marsaglia (1972) is then employed to generate points from a distribution on the unit 3-sphere. Using n such random points, the statistics Y_n^{odd} , Y_n^{even} and Y_n defined in Section 2 are computed. With Monte-Carlo samples of size 5000 for $n = 100$, and 20000 for $n = 5, 10, 20$ and 40, the percent points of the above statistics are estimated for the significance levels $\alpha = .20, .10, .05, .025, .01$ and $.001$. Those empirical finite sample percentage points are further smoothed in the following manner. For each level α and sample size n , let a smooth function $\hat{Y}_{\alpha,n}$ be defined by

$$\hat{Y}_{\alpha,n} = a_{\alpha,\infty} + \frac{4}{n^2} (a_{\alpha,2} - a_{\alpha,\infty}) + b_{\alpha,n} \left(\frac{1}{n} - \frac{2}{n^2} \right) \quad (3.1)$$

Here $\hat{Y}_{\alpha,n}$ is the estimated Y value for each n at level α . $a_{\alpha,2}$ and $a_{\alpha,\infty}$ are the Y values at level α for $n = 2$ and $n = \infty$, respectively. $a_{\alpha,2}$ is found from Theorem 2, and $a_{\alpha,\infty}$ from the Laguerre transform method to be described in Section 4. $b_{\alpha,n}$ is the estimated slope obtained from the original Monte-Carlo results via the straightforward linear regression. It should be noted that the smooth function (3.1) coincides with the known values of Y when $n = 2$ or $n = \infty$.

The estimated smoothed percentage points of the three statistics are presented in Table 6.1 of Section 6.

§4. Evaluation of the limiting distributions via the Laguerre transform

We have seen in Section 2 that Beran's and Giné's limiting statistics under the null hypothesis of uniformity are infinite sums of independent and scaled Chi-square variates, i.e., Y^{odd} , Y^{even} and Y as given in (2.6), (2.7) and (2.8). In this section, we discuss numerical evaluation of Y^{odd} , Y^{even} and Y and propose a Laguerre transform approach to be described.

It is natural to decompose the infinite sum $Y = \sum_{j=1}^{\infty} a_j^2 z_j$ in (2.8) into two parts, the sum of the first N variates and the remainder, i.e.,

$$Y = S_N + V_N ; \quad S_N = \sum_{j=1}^N a_j^2 z_j ; \quad V_N = \sum_{j=N+1}^{\infty} a_j^2 z_j . \quad (4.1)$$

Correspondingly, let

$$Y^{\text{odd}} = S_N^{\text{odd}} + V_N^{\text{odd}} ; \quad Y^{\text{even}} = S_N^{\text{even}} + V_N^{\text{even}} \quad (4.2)$$

where S_N^{odd} and V_N^{odd} are the sums of odd index terms of S_N and V_N , respectively. The variates S_N^{even} and V_N^{even} are defined similarly for even index terms. S_N is the finite sum of independent scaled Chi-square variates and such linear combinations may be regarded as positive-definite quadratic forms in normal variables. Many papers have been published on the distribution of such quadratic forms, and the reader is referred to Johnson and Kotz (1970, Ch. 29) for a comprehensive survey of the literature. An excellent approach to the numerical evaluation of such distributions is that of Johnson, Kotz and Boyd (1967). They expand the distribution function in a series of generalized Laguerre functions and evaluate it efficiently by taking advantage of the recurrence relation of the Laguerre functions. When this procedure is applied directly to the distribution of

S_N , however, one encounters numerical difficulty. The coefficients a_j^2 decreases rapidly and the distributions of $a_j^2 z_j$ become very concentrated, resulting in quite slow convergence of the corresponding Laguerre series. Our procedure restructures the method of Johnson, Kotz and Boyd and overcomes this numerical difficulty.

The Laguerre transform method for convolving functions has been introduced by Keilson and Nunn (1979), Keilson, Nunn and Sumita (1981), and further studied by Sumita (1981). The method has advantages of accuracy and speed which make it an attractive candidate for problems of this type. The Laguerre transform method has peculiarities and limitations, however, which require careful refinement for particular contexts, such as that here. The reader is referred to three basic papers for the underlying theory. A brief summary is given in Appendix A for the convenience of the reader. Our basic strategy is to approximate Y^{odd} and Y^{even} by

$$Y_N^{\text{odd}*} = S_N^{\text{odd}} + V_N^{\text{odd}*} ; \quad Y_N^{\text{even}*} = S_N^{\text{even}} + V_N^{\text{even}*} \quad (4.3)$$

where $V_N^{\text{odd}*}$ and $V_N^{\text{even}*}$ are the Gamma variates having the same first two moments of V_N^{odd} and V_N^{even} , respectively. Correspondingly, Y is approximated by

$$Y_N^* = Y_N^{\text{odd}*} + Y_N^{\text{even}*} = S_N + V_N^* \quad (4.4)$$

In the subsections to follow, the Laguerre transform procedure is described for evaluating the distribution of $Y_N^{\text{odd}*}$, $Y_N^{\text{even}*}$ and Y_N^* . The

validity of the results is also examined.

(A) The Laguerre sharp coefficients of the Gamma variate $\Gamma(\alpha, 2\beta)$

Let $\Gamma(\alpha, 2\beta)$, $\alpha, \beta > 0$, be the Gamma variate with p.d.f.

$$g(x) = \frac{1}{\Gamma(\alpha)(2\beta)^\alpha} x^{\alpha-1} e^{-\frac{1}{2\beta}x}, \quad 0 \leq x < \infty. \quad (4.5)$$

It is clear that the variates $a_j^2 Z_j$ belong to this family with $\alpha = j + \frac{1}{2}$ and $\beta = a_j^2$. Hence the Laguerre sharp coefficients $(g_n^\#)_0^\infty$ of $g(x)$ provides a basic tool for the procedure. From the Laplace transform $\gamma(s) = \int_0^\infty e^{-sx} g(x) dx = (1 + 2\beta s)^{-\alpha}$ and the identity $T_g^\#(u) = \sum_{n=0}^\infty g_n^\# u^n = \gamma(\frac{1}{2} \cdot \frac{1+u}{1-u})$, those coefficients are found by (cf. Sumita (1981), Section 6.2)

$$g_n^\# = (1+\beta)^{-\alpha} \sum_{m=0}^n b_{n-m} c_m \quad (4.6)$$

where

$$b_n = \prod_{r=1}^n (1 - \frac{1+\alpha}{r}), \quad n \geq 1, \quad b_0 = 1$$

$$c_n = (\frac{1-\beta}{1+\beta})^n \prod_{r=1}^n (1 - \frac{1-\alpha}{r}), \quad n \geq 1, \quad c_0 = 1. \quad (4.7)$$

The accuracy and efficiency of the Laguerre transform method depend heavily on one's ability to represent the functions present with a sequence of Laguerre coefficients of reasonable length (say, at most around 500 coefficients to attain 5 digits accuracy). As studied theoretically in Keilson, Nunn and Sumita (1981) and Sumita (1981), the Laguerre transform method, when applied in a straightforward manner, cannot tolerate functions

too closely concentrated at zero or functions too great in extent.

For the Gamma variate, this point can be observed explicitly in (4.6)

and (4.7). When α is extremely large, b_n and c_n become so large in

absolute value that the computer may not tolerate them. The other

numerical difficulty arises when β is extremely small or large. In

this case, the ratio $\left| \frac{1-\beta}{1+\beta} \right|$ becomes very close to 1, and one would

expect (c_n) to have a long tail. This, in turn, implies a long tail

of $g_n^\#$. Fortunately, these numerical difficulties can be avoided by

taking advantage of the divisibility of the Gamma variates and employ-

ing scaling. In brief, the first difficulty can be solved through the

identity $\gamma(s) = [\gamma_M(s)]^M$, where $M > 0$ and $\gamma_M(s) = (1 + 2\beta s)^{-\alpha/M}$ corre-

sponding to the Gamma variate $\Gamma(\alpha/M, 2\beta)$. By an appropriate choice of

positive integer M , the Laguerre sharp coefficients of $\Gamma(\alpha/M, 2\beta)$ are

obtained with reasonable length. We then convolve them M times on the

lattice to recover the Laguerre sharp coefficients of the original

Gamma variate $\Gamma(\alpha, 2\beta)$. For the second numerical difficulty, we replace

β by $c\beta$ so that the ratio $\left| \frac{1-c\beta}{1+c\beta} \right|$ becomes well below 1. Then (c_n) in

(4.7) decreases rapidly in absolute value and therefore $g_n^\#$ decreases

rapidly. After the inversion of the Laguerre sharp representation, the

proper scale factor for the probability density function is restored.

(B) Algorithms for finding the sharp coefficients of $Y^{\text{odd}*}$, $Y^{\text{even}*}$ and Y^*

(1) For a desired small variance of V_N , select N for S_N and V_N .

Then choose the dividing factor M for V_N , and the scale factor c in keeping

with the conditions of (A). The number L of Laguerre coefficients before

truncation will be discussed in (C) below.

(2) Using (4.6) and (4.7) with $\alpha = j + \frac{1}{2}$ and $\beta = ca_j^2$, calculate $(g(j)_n)_0^L$ of $ca_j^2 z_j$ for $1 \leq j \leq N$.

(3) Obtain $(g_{N,n}^{\text{odd}})_0^L$ of cS_n^{odd} by convolving $(g(j)_n)_0^L$ for j odd, $1 \leq j \leq N$. Obtain $(g_{N,n}^{\text{even}})_0^L$ of cS_n^{even} similarly.

(4) Calculate the means μ_{VN}^{odd} and μ_{VN}^{even} and the variances $\sigma_{VN}^{\text{odd}2}$ and $\sigma_{VN}^{\text{even}2}$ of V_N^{odd} and V_N^{even} , respectively, from (2.6), (2.7) and Proposition

3. Using (4.6) and (4.7) with $\alpha = (\mu_{VN}^{\text{odd}}/\sigma_{VN}^{\text{odd}})^2/M$ and $\beta = c \cdot \sigma_{VN}^{\text{odd}2}/2\mu_{VN}^{\text{odd}}$, calculate $(h^{\text{odd}\#}(M,N)_n)_0^L$. By convolving $(h^{\text{odd}\#}(M,N)_n)_0^L$ M times with itself, find $(h_{N,n}^{\text{odd}\#})_0^L$ of $cV_N^{\text{odd}*}$. Obtain $(h_{N,n}^{\text{even}\#})_0^L$ of $cV_N^{\text{even}*}$ similarly.

(5) Calculate $(f_{N,n}^{\text{odd}\#})_0^L$ of $cY_N^{\text{odd}*}$ by convolving $(g_{N,n}^{\text{odd}\#})_0^L$ and $(h_{N,n}^{\text{odd}\#})_0^L$. Obtain $(f_{N,n}^{\text{even}\#})_0^L$ of $cY_N^{\text{even}*}$ similarly.

(6) Finally, calculate $(f_{N,n}^{\#})_0^L$ of cY_N^* by convolving $(f_{N,n}^{\text{odd}\#})_0^L$ and $(f_{N,n}^{\text{even}\#})_0^L$.

Remark

When one has the Laguerre sharp coefficients $(f_n^{\#})_0^{\infty}$ of a p.d.f. $f(x)$ on $(0, \infty)$, the inversion of $(f_n^{\#})_0^{\infty}$ to the values of $f(x)$ and its survival function $\bar{F}(x) = \int_x^{\infty} f(x)dx$ can be done in the following manner (cf. Keilson and Nunn (1980)).

(R1) Calculate $(\ell_n(x))_0^{\infty}$ for $0 \leq x < \infty$ by

$$\ell_{n+1}(x) = \frac{1}{n+1} [(2n+1-x)\ell_n(x) - n\ell_{n-1}(x)] \quad , \quad n \geq 1$$

$$\text{where } \ell_0(x) = e^{-\frac{1}{2}x}.$$

(R2) Obtain $f_n^+ = \sum_{m=0}^n f_m^{\#}$ and calculate $f(x) = \sum_{m=0}^{\infty} f_n^+ \ell_n(x)$.

(R3) Calculate $\bar{f}_n^+ = -2 \sum_{m=0}^{\infty} (-1)^m f_{n+1+m}^{\#}$ and $\bar{F}(x) = \sum_{n=0}^{\infty} \bar{f}_n^+ \ell_n(x)$.

We note that the Laguerre transform bypasses numerical integration. We also note that the algorithm with (R1), (R2) and (R3) produces $\bar{F}_N^*(x/c)$ where $\bar{F}_N^*(x) = P[Y_N^* > x]$.

(C) Validation of the results

For the calculation of the distributions needed, the scaling factor is taken to be $c = 40$. The two values $N = 2$ and $N = 4$ are used and the dividing factor $M = 12$ is chosen for V_2^* and $M = 25$ for V_4^* . The length of the Laguerre sharp coefficients is $L = 502$, which provides 12 digits accuracy of the p.d.f. of $a_j^2 z_j$ for $1 \leq j \leq 4$.

There are two different factors which introduce numerical errors, the truncation of the Laguerre sharp coefficients and the Gamma approximation of V_N . In general, it is quite hard to quantify truncation error of the Laguerre coefficients. (Such error bounding has its counterpart in Fourier series theory, where error bounding is known to be extremely difficult.) Theoretical error bounds are available, so far, only for a certain family of functions (cf. Keilson and Sumita (1981) and Sumita (1981)). Extensive numerical evidence, however, suggests that when one chooses L large enough to attain a given accuracy for the following identities, then the function values are likely to satisfy the same accuracy.

$$f(0+) = - \sum_{n=0}^{\infty} n f_n^{\#} \quad (4.8)$$

$$\int_0^{\infty} f(x) dx = \sum_{n=0}^{\infty} (-1)^n f_n^{\#} \quad (4.9)$$

$$\int_0^{\infty} x f(x) dx = 4x \sum_{n=0}^{\infty} (-1)^n n f_n^{\#} \quad (4.10)$$

$$\int_0^{\infty} x^2 d(x) dx = 16 \times \sum_{n=0}^{\infty} (-1)^n n^2 f_n^{\#} . \quad (4.11)$$

For $N = 4$, $(f_{4,n}^{\#})_0^{502}$ of Y_4^* provides 10 digits accuracy for all equations (4.8) through (4.10) and the truncation error seems to be negligible.

Even though we may expect the Gamma approximations for the remainders V_N^{odd} , V_N^{even} , and V_N to introduce little error since $\text{Var}[V_N]$ drops rapidly (e.g., $\text{Var}[V_4] = 6.6 \times 10^{-4}$), no analytical justification is available, and we are forced to take indirect means for testing the validity of the approximation. To test this validity we note that $Y_N^* \xrightarrow{d} Y$ as $N \rightarrow \infty$. The c.d.f.'s of Y_2^* and Y_4^* are calculated and compared. The absolute difference of the two c.d.f.'s is found numerically to be bounded by 1×10^{-5} . For a second accuracy check of the Gamma approximation, the third and the fourth cumulants of $Y_4^{\text{odd}*}$ and $Y_4^{\text{even}*}$ are calculated using the Laguerre sharp coefficients and compared with exact values. The higher moments of $Y_4^{\text{odd}*}$ and $Y_4^{\text{even}*}$ can be found from (4.10), (4.11) and (R3) of (B) with the identity $f_n^{\#} = f_n^+ - f_{n-1}^+$, $n \geq 1$, and $f_n^{\#} = f_0^+$. The cumulants needed are then obtained by

$$K_3 = \mu_3 - 3\mu_1\mu_2 + 2\mu_1^3 \quad (4.12)$$

$$K_4 = \mu_4 - 3\mu_2^2 - 4\mu_1\mu_3 + 12\mu_1^2\mu_2 - 6\mu_1^4$$

where $\mu_i = E[X^i]$. The corresponding exact values are given in Proposition 3. The computations are shown in Table 4.1. The absolute difference is bounded by 3×10^{-6} .

Finally, we study the asymptotic behavior of the survival function of Y , and compare it with the values computed via the Laguerre transform method. From the Laplace transform of Y , one finds through asymptotic analysis that

$$\bar{F}_Y(x) = P[Y > x] \sim \frac{e^{-2x}}{\sqrt{\pi}} \left[2\lambda_0 \sqrt{2x} + \left(\frac{\lambda_0}{\sqrt{2}} + \lambda_1 \sqrt{2} \right) \frac{1}{\sqrt{x}} \right] \quad (4.13)$$

as $x \rightarrow +\infty$, where $\lambda_0 \approx 4.95221$ and $\lambda_1 \approx -4.27140$. Details are given in Appendix B. In Figure 4.2, the asymptotic expansion of $\bar{F}_Y(x)$ in (4.13) is plotted with the survival function of Y_4^* derived via the Laguerre transform method. Their absolute difference is found to be bounded by 1×10^{-5} for $x > 3.5$.

A similar approximation for the remainder of an infinite sum of independent variates is employed in Sumita (1979) to evaluate the multiple convolutions of the Logistic variates. There a direct accuracy check is possible and accuracy to seven decimal places is attained.

All calculations were carried out on a DEC10 computer, in a time-sharing mode using APL as the programming language. Relevant formulae are coded in a straightforward manner, with no attempt made to optimize the subroutines for speed and accuracy. In spite of this, the results displayed here were obtained with CPU times in seconds with no evidence of numerical problems.

$K_3 (Y^{\text{odd}})$	0.3752193777	$K_4 (Y^{\text{odd}})$	0.5625201543
$K_3 (Y_4^{\text{odd}*})$	0.3752164153	$K_4 (Y_4^{\text{odd}*})$	0.5625200535
Absolute Difference	0.0000029624	Absolute Difference	0.0000001008
$K_3 (Y^{\text{even}})$	0.0098016768	$K_4 (Y^{\text{even}})$	0.0036637420
$K_3 (Y_4^{\text{even}*})$	0.0098007718	$K_4 (Y_4^{\text{even}*})$	0.0036637195
Absolute Difference	0.0000009050	Absolute Difference	0.0000000225

Table 4.1. The cumulants: Y^{odd} and Y^{even} vs. $Y_4^{\text{odd}*}$ and $Y_4^{\text{even}*}$

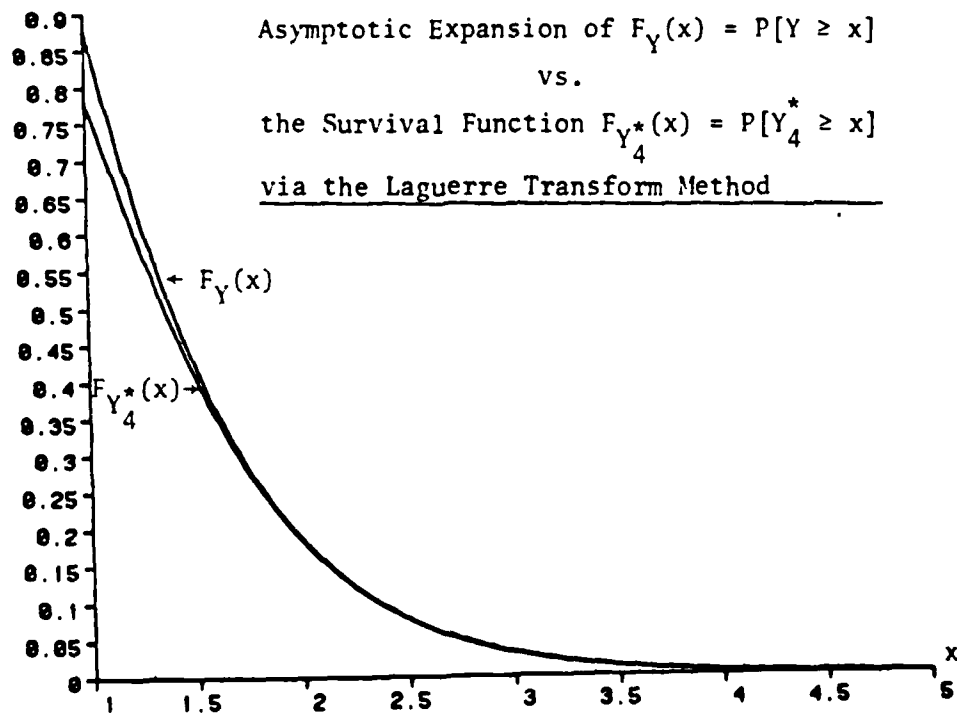


Figure 4.2

§5. Application

In an experiment to determine the effect of different visual stimuli, six cats were divided into three groups of two, and each group was subjected to a different visual stimulus: horizontally polarized light (H), vertically polarized light (V), and unpolarized or "normal" light (N). Orientations of the dendritic fields were then measured at 15 to 16 sites in the retinas of each of the six cats. The complete data set, in coordinates described by Figure 5.1 below, is given in Appendix C.

The orientation of the dendritic fields of the two cats exposed to normal light is, presumably, uniform, with no preferred orientation, and the question is: what is the effect of polarized light? Hence we wish to test the null hypothesis of uniformity of orientation (on the unit hemisphere) of the dendritic fields of the H and V groups. To test this hypothesis, we use Giné's statistic $\gamma_n^{\text{even}} = n^{-1} \sum_{i=1}^n \sum_{j=1}^n (\frac{1}{2} - (\frac{2}{\pi}) \sin \hat{X}_i X_j)$, where $\hat{X}_i X_j$ is the angle between the observations X_i and X_j . Since X_i is represented by a vector $\underline{x}_i = [\cos\phi_i \sin\theta_i, \sin\phi_i \sin\theta_i, \cos\theta_i]$ with $||\underline{x}_i||_2 = 1$, one easily finds that $\cos \hat{X}_i X_j = \underline{x}_i^T \underline{x}_j$ and therefore

$$(5.1) \quad \hat{X}_i X_j = \arccos \left[\frac{1}{2} \cos(\theta_i - \theta_j) \{1 + \cos(\phi_i - \phi_j)\} \right. \\ \left. + \frac{1}{2} \cos(\theta_i + \theta_j) \{1 - \cos(\phi_i - \phi_j)\} \right] .$$

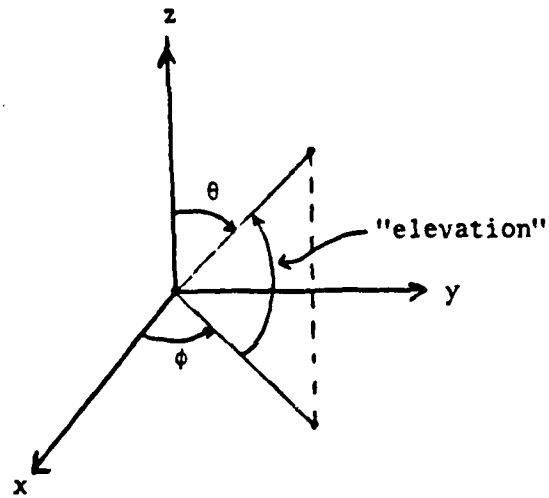


Figure 5.1

A summary table of the P-values using Table 6.1 is given below.

Cats	n	y_n^{even}	P-value	Cats	n	y_n^{even}	P-value
H_1	15	.4508	>> .20	$H_1 + H_2$	31	.9181	\approx .040
H_2	16	.7321	\approx .120				
V_1	16	.4923	>> .20	$V_1 + V_2$	32	.8603	\approx .056
V_2	16	.6239	> .20				
N_1	15	.6260	> .20	$N_1 + N_2$	31	.4421	>> .20
N_2	16	.6076	> .20				

Table 5.1

As the table shows, the test of uniformity is not significant for any individual cat. Grouping H_1 and H_2 or V_1 and V_2 , however, does give significant results, whereas grouping N_1 and N_2 is still not significant.

§6. Graphs and Tables

In Table 6.1, the estimated smoothed percentage points for various values of n and significance levels α are presented, using the method described in Section 3. For $n = 2$ and $n = \infty$, the values are taken from the formulae in Theorem 2 and Table 6.2, respectively. For comparison purposes, we record the approximate significance points of Prentice (1978). The quantiles of the limiting distributions of Y^{odd} , Y^{even} and Y are given in Table 6.2, where $\alpha = P[Y > y]$, etc. More detailed values of the distributions are presented in Table 6.3, with step size 0.05. Figure 6.1 and Figure 6.2 show the graphs of the c.d.f. and the p.d.f. of Y^{odd} , Y^{even} and Y . All the values associated with the limiting distributions are calculated by the Laguerre transform method described in Section 4.

Table 6.1

Selected Smoothed Percentage Points for Y_n^{odd} , Y_n^{even} and Y_n

Sample Size	Statistics	Significance level α					
		.20	.10	.05	.025	.01	.001
2 (Exact)	Y_n^{odd}	1.4097	1.5903	1.7129	1.7978	1.8725	1.9597
	Y_n^{even}	0.6180	0.7225	0.8012	0.8585	0.9102	0.9715
	Y_n	1.9004	2.2084	2.4359	2.5991	2.7458	2.9195
5	Y_n^{odd}	1.4171	1.7788	2.0771	2.3456	2.6704	3.3946
	Y_n^{even}	0.6269	0.7345	0.8348	0.9323	1.0598	1.3838
	Y_n	1.9132	2.2767	2.6154	2.9387	3.3338	4.3277
10	Y_n^{odd}	1.4162	1.8063	2.1534	2.4802	2.8914	3.8606
	Y_n^{even}	0.6353	0.7490	0.8567	0.9619	1.1000	1.4476
	Y_n	1.9278	2.3124	2.6803	3.0389	3.4927	4.6366
20	Y_n^{odd}	1.4151	1.8135	2.1832	2.5385	2.9935	4.0913
	Y_n^{even}	0.6405	0.7582	0.8696	0.9777	1.1183	1.4657
	Y_n	1.9372	2.3327	2.7137	3.0866	3.5652	4.7610
40	Y_n^{odd}	1.4144	1.8155	2.1959	2.5654	3.0424	4.2060
	Y_n^{even}	0.6434	0.7634	0.8765	0.9858	1.1270	1.4712
	Y_n	1.9423	2.3434	2.7306	3.1098	3.5997	4.8156
100	Y_n^{odd}	1.4140	1.8161	2.2029	2.5808	3.0712	4.2747
	Y_n^{even}	0.6452	0.7666	0.8809	0.9908	1.1320	1.4735
	Y_n	1.9455	2.3501	2.7409	3.1236	3.6199	4.8460
∞	Y_n^{odd}	1.41363	1.81631	2.20727	2.59079	3.09000	4.32040
	Y_n^{even}	0.64643	0.76879	0.88384	0.99413	1.13534	1.47452
	Y	1.94776	2.35459	2.74772	3.13268	3.63309	4.86522
<u>Prentice (1978)</u>							
(+4)	Y_n^{odd}	-	1.7750	2.1750	-	3.0500	4.3000
(+16)	Y_n^{even}	-	0.7625	0.8813	-	1.1313	1.4680

Table 6.2. The Quantiles of the Limiting Distributions

α	$Y(\text{ODD})$	$Y(\text{EVEN})$	Y
0.999	0.18563	0.16005	0.47887
0.995	0.22343	0.18567	0.55003
0.990	0.24414	0.20184	0.58984
0.975	0.28309	0.22580	0.65944
0.950	0.32650	0.24992	0.73038
0.900	0.39123	0.28234	0.82748
0.850	0.44743	0.30846	0.90397
0.800	0.50073	0.33171	0.97212
0.750	0.55346	0.35351	1.03633
0.700	0.60686	0.37460	1.09897
0.650	0.66181	0.39553	1.16151
0.600	0.71904	0.41674	1.22509
0.550	0.77939	0.43850	1.29091
0.500	0.84376	0.46113	1.35992
0.450	0.91324	0.48501	1.43340
0.400	0.98921	0.51060	1.51289
0.350	1.07352	0.53849	1.60028
0.300	1.16916	0.56940	1.69861
0.250	1.28015	0.60472	1.81215
0.200	1.41363	0.64643	1.94776
0.150	1.58259	0.69823	2.11887
0.100	1.81631	0.76879	2.35459
0.050	2.20727	0.88384	2.74772
0.025	2.59079	0.99413	3.13268
0.010	3.09000	1.13534	3.63309
0.005	3.46337	1.23944	4.00709
0.001	4.32040	1.47452	4.86522

(e.g., $P[Y > 0.47887] = 0.999$)

Table 6.3. The Cumulative Distribution Functions of Y^{odd} , Y^{even} and Y

x	$F\text{-ODD}(x)$	$F\text{-EVEN}(x)$	$F(x)$	x	$F\text{-ODD}(x)$	$F\text{-EVEN}(x)$	$F(x)$
0.00	0.00000	0.00000	0.00000	2.55	0.97308	1.00000	0.92899
0.05	0.00000	0.00000	0.00000	2.60	0.97542	1.00000	0.93500
0.10	0.00000	0.00000	0.00000	2.65	0.97756	1.00000	0.94051
0.15	0.00005	0.00025	0.00000	2.70	0.97952	1.00000	0.94557
0.20	0.00180	0.00887	0.00000	2.75	0.98131	1.00000	0.95021
0.25	0.01147	0.05008	0.00000	2.80	0.98295	1.00000	0.95446
0.30	0.03353	0.13223	0.00000	2.85	0.98444	1.00000	0.95836
0.35	0.06670	0.24167	0.00000	2.90	0.98581	1.00000	0.96194
0.40	0.10739	0.36068	0.00006	2.95	0.98706	1.00000	0.96522
0.45	0.15235	0.47607	0.00040	3.00	0.98820	1.00000	0.96822
0.50	0.19930	0.58021	0.00170	3.05	0.98924	1.00000	0.97096
0.55	0.24673	0.66972	0.00500	3.10	0.99019	1.00000	0.97348
0.60	0.29367	0.74401	0.01152	3.15	0.99105	1.00000	0.97578
0.65	0.33945	0.80405	0.02228	3.20	0.99184	1.00000	0.97788
0.70	0.38368	0.85158	0.03792	3.25	0.99257	1.00000	0.97981
0.75	0.42606	0.88860	0.05860	3.30	0.99323	1.00000	0.98157
0.80	0.46643	0.91703	0.08408	3.35	0.99383	1.00000	0.98318
0.85	0.50471	0.93863	0.11185	3.40	0.99438	1.00000	0.98465
0.90	0.54085	0.95487	0.14718	3.45	0.99488	1.00000	0.98600
0.95	0.57486	0.96699	0.18331	3.50	0.99533	1.00000	0.98723
1.00	0.60678	0.97597	0.22144	3.55	0.99575	1.00000	0.98835
1.05	0.63665	0.98258	0.26083	3.60	0.99613	1.00000	0.98937
1.10	0.66455	0.98743	0.30083	3.65	0.99647	1.00000	0.99031
1.15	0.69056	0.99095	0.34085	3.70	0.99679	1.00000	0.99116
1.20	0.71476	0.99351	0.38043	3.75	0.99707	1.00000	0.99194
1.25	0.73726	0.99536	0.41917	3.80	0.99734	1.00000	0.99266
1.30	0.75813	0.99669	0.45670	3.85	0.99757	1.00000	0.99331
1.35	0.77748	0.99764	0.49303	3.90	0.99779	1.00000	0.99390
1.40	0.79539	0.99833	0.52776	3.95	0.99799	1.00000	0.99444
1.45	0.81196	0.99882	0.56085	4.00	0.99817	1.00000	0.99493
1.50	0.82727	0.99916	0.59224	4.05	0.99833	1.00000	0.99538
1.55	0.84141	0.99941	0.62191	4.10	0.99848	1.00000	0.99580
1.60	0.85445	0.99958	0.64985	4.15	0.99862	1.00000	0.99617
1.65	0.86647	0.99971	0.67609	4.20	0.99874	1.00000	0.99651
1.70	0.87755	0.99979	0.70067	4.25	0.99886	1.00000	0.99682
1.75	0.88774	0.99986	0.72364	4.30	0.99896	1.00000	0.99711
1.80	0.89713	0.99990	0.74506	4.35	0.99905	1.00000	0.99737
1.85	0.90576	0.99993	0.76501	4.40	0.99914	1.00000	0.99760
1.90	0.91369	0.99995	0.78355	4.45	0.99922	1.00000	0.99782
1.95	0.92097	0.99997	0.80076	4.50	0.99929	1.00000	0.99801
2.00	0.92767	0.99998	0.81671	4.55	0.99935	1.00000	0.99819
2.05	0.93381	0.99998	0.83148	4.60	0.99941	1.00000	0.99835
2.10	0.93944	0.99999	0.84514	4.65	0.99946	1.00000	0.99850
2.15	0.94461	0.99999	0.85777	4.70	0.99951	1.00000	0.99863
2.20	0.94935	0.99999	0.86942	4.75	0.99956	1.00000	0.99876
2.25	0.95370	1.00000	0.88017	4.80	0.99960	1.00000	0.99887
2.30	0.95768	1.00000	0.89000	4.85	0.99963	1.00000	0.99897
2.35	0.96133	1.00000	0.89921	4.90	0.99967	1.00000	0.99906
2.40	0.96467	1.00000	0.90762	4.95	0.99970	1.00000	0.99915
2.45	0.96772	1.00000	0.91535	5.00	0.99972	1.00000	0.99923
2.50	0.97052	1.00000	0.92246				

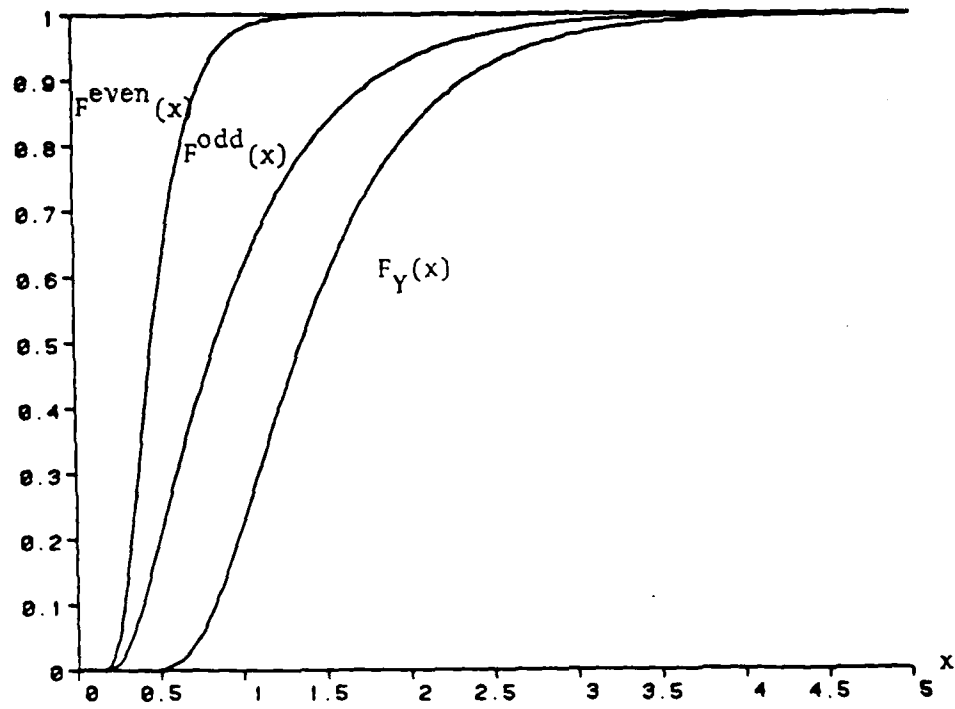


Fig. 6.1. The cumulative distribution functions of Y^{odd} , Y^{even} and Y

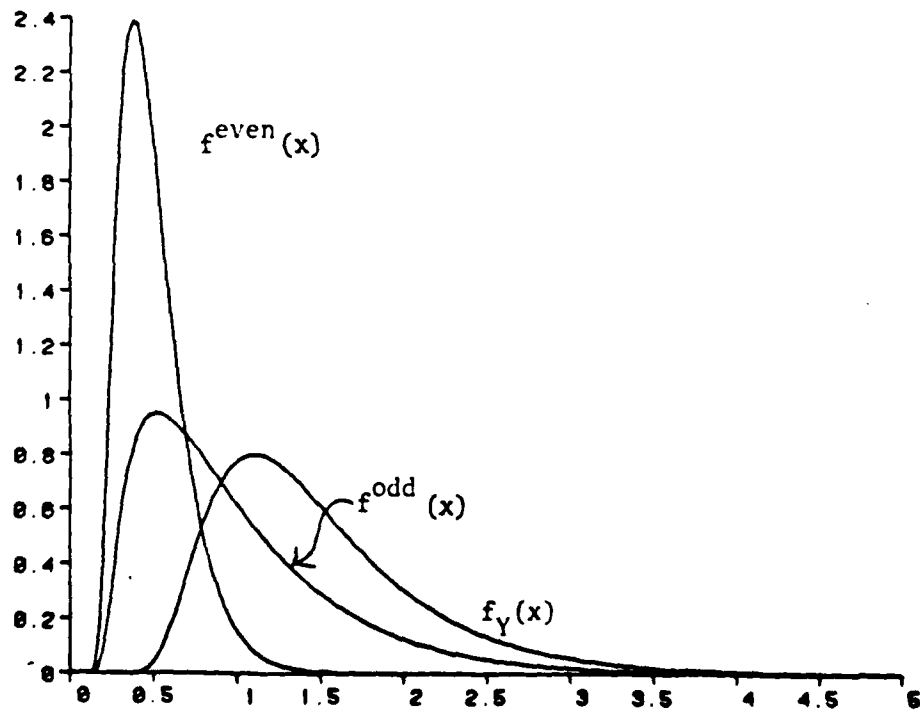


Fig. 6.2. The probability density functions of Y^{odd} , Y^{even} and Y

APPENDIX A

The Laguerre Transform

The Laguerre polynomials $L_n(x)$, defined by the Rodrigues formula $L_n(x) = \frac{1}{n!} e^x \left(\frac{d}{dx}\right)^n (x^n e^{-x})$, form a set of orthonormal polynomials with weighting function $w(x) = e^{-x}$ on $(0, \infty)$ (see, e.g., Szego (1975)). The associated Laguerre functions $\ell_n(x) = e^{-\frac{1}{2}x} L_n(x)$ then provide an orthonormal basis in $L_2(0, \infty)$. For any $f(x) \in L_2(0, \infty)$, one has the Fourier-Laguerre expansion

$$(A.1) \quad f(x) = \sum_{n=0}^{\infty} f_n^+ \ell_n(x) ; \quad f_n^+ = \int_0^{\infty} f(x) \ell_n(x) dx .$$

Let $T_f^+(u) = \sum_{n=0}^{\infty} f_n^+ u^n$ and $T_f^\#(u) = \sum_{n=0}^{\infty} f_n^\# u^n \stackrel{\text{def}}{=} (1-u) T_f^+(u)$. Since $\sum_{n=0}^{\infty} \ell_n(x) u^n = (1-u)^{-1} \exp\{-\frac{1}{2}x(1+u)(1-u)^{-1}\}$, one has, when $f(x)$ is integrable on $(0, \infty)$,

$$\begin{aligned} T_f^\#(u) &= (1-u) \sum_{n=0}^{\infty} f_n^+ u^n = (1-u) \sum_{n=0}^{\infty} u^n \int_0^{\infty} f(x) \ell_n(x) dx \\ &= \int_0^{\infty} f(x) \exp\{-\frac{1}{2}x(1+u)(1-u)^{-1}\} dx , \quad 0 \leq u < 1 , \end{aligned}$$

i.e.,

$$(A.2) \quad T_f^\#(u) = \phi\left(\frac{1+u}{2(1-u)}\right)$$

where $\phi(s) = \int_0^{\infty} e^{-sx} f(x) dx$ is the Laplace transform of $f(x)$. Let $f(x) * g(x) = \int_0^x f(x-y) g(y) dy$. Since $\phi_{f*g}(s) = \phi_f(s) \phi_g(s)$,

$$(A.3) \quad T_{f*g}^{\#}(u) = T_f^{\#}(u)T_g^{\#}(u)$$

for functions $f(x)$, $g(x)$ that are both in $L_1(0, \infty)$ and $L_2(0, \infty)$. From (A.3) one obtains

$$(A.4) \quad (f*g)_n^{\#} = \sum_{m=0}^n f_{n-m}^{\#} g_m^{\#} .$$

The transformation via (A.3) maps functions $f(x)$, $g(x)$ into sequences $(f_n^{\#})(g_n^{\#})$, and their continuum convolution $f(x)*g(x)$ is mapped into a lattice convolution and then back onto the continuum via $f_n^{\dagger} = \sum_{m=0}^n f_m^{\#}$ and the representation (A.1). This transformation procedure was introduced originally in Keilson and Nunn (1980). The Laguerre transform was extended subsequently in Keilson, Nunn and Sumita (1981) to handle functions on the full continuum, and further studied by Sumita (1981).

APPENDIX B

Asymptotic Behavior of $\bar{F}_Y(x) = P[Y > x]$ as $x \rightarrow +\infty$

In this appendix, we derive the asymptotic expression of $\bar{F}_Y(x)$ as $x \rightarrow +\infty$. The Laplace transform of the p.d.f. of Y is given by

$$(B.1) \quad \phi_Y(s) = \prod_{j=1}^{\infty} \left(\frac{\theta_j}{\theta_j + s} \right)^{j+1/2}; \quad \theta_j = \frac{1}{2a_j^2}.$$

Let

$$(B.2) \quad \alpha(s) = \left(\frac{\theta_1}{\theta_1 + s} \right)^{3/2}; \quad \lambda(s) = \prod_{j=2}^{\infty} \left(\frac{\theta_j}{\theta_j + s} \right)^{j+1/2}.$$

We note that $\lambda(s)$ is regular for $\text{Re}(s) > -\theta_2$. Let $\bar{F}_Y(x) = P[Y > x]$ and define $\psi(w) = L\{e^{\theta_1 x} \bar{F}_Y(x)\} = \int_0^{\infty} e^{-(w - \theta_1)x} \bar{F}_Y(x) dx$. Then, from (B.1) and (B.2), one has

$$(B.3) \quad \psi(w) = \frac{1 - \phi_Y(w - \theta_1)}{w - \theta_1} = \frac{1 - \alpha(w - \theta_1)}{w - \theta_1} \cdot \lambda(w - \theta_1) + \frac{1 - \lambda(w - \theta_1)}{w - \theta_1}.$$

After a little algebra, one finds, from (B.2), that

$$(B.4) \quad \frac{1 - \alpha(w - \theta_1)}{w - \theta_1} = \frac{\sqrt{\theta_1}}{w^{3/2}} + \frac{1}{\sqrt{\theta_1} w} + \sum_{n=3}^{\infty} (-1)^n \left(\frac{1}{\theta} \right)^{\frac{n-1}{2}} \frac{1}{w^{\frac{n-3}{2}}}.$$

On the other hand, $\lambda(w - \theta_1)$ and $\frac{1 - \lambda(w - \theta_1)}{w - \theta_1}$ are regular for $\text{Re}(w) > \theta_1 - \theta_2$ with $\theta_1 - \theta_2 < 0$. Hence, one has the Taylor expansion

$$(B.5) \quad \lambda(w - \theta_1) = \sum_{n=0}^{\infty} \lambda_n w^n.$$

From (B.3), (B.4), (B.5) and the regularity of $\frac{1 - \lambda(w - \theta_1)}{w - \theta_1}$ for $\text{Re}(w) > \theta_1 - \epsilon_2$, we obtain the asymptotic expression of $\psi(w)$ near zero, i.e.,

$$(B.6) \quad \psi(w) \sim \frac{\lambda_0 \sqrt{\theta_1}}{w^{3/2}} + \left(\frac{\lambda_0}{\sqrt{\theta_1}} + \lambda_1 \sqrt{\theta_1} \right) \frac{1}{\sqrt{w}} \quad \text{as } w \rightarrow 0+.$$

This then implies that (see, e.g., Widder (1946), p. 192)

$$^{\dagger}(B.7) \quad \bar{F}_Y(x) \sim \frac{e^{-2x}}{\sqrt{\pi}} [2\lambda_0 \sqrt{2x} + \left(\frac{\lambda_0}{\sqrt{2}} + \lambda_1 \sqrt{2} \right) \frac{1}{\sqrt{x}}] \quad \text{as } x \rightarrow +\infty,$$

where $\theta_1 = 2$ is substituted. The constants λ_0 and λ_1 are obtained from

$$\lambda(w - \theta_1) = \prod_{j=2}^{\infty} \left(\frac{\theta_j}{\theta_j - \theta_1 + w} \right)^{j+1/2} = \sum_{n=0}^{\infty} \lambda_n w^n, \text{ i.e.,}$$

$$(B.8) \quad \begin{cases} \lambda_0 = \prod_{j=2}^{\infty} \left(\frac{\theta_j}{\theta_j - \theta_1} \right)^{j+1/2} \approx 4.952213394 \\ \lambda_1 = - \sum_{j=2}^{\infty} \frac{j+1/2}{\theta_j - \theta_1} \prod_{\substack{r=2 \\ r \neq j}}^{\infty} \left(\frac{\theta_r}{\theta_r - \theta_1} \right)^{r+1/2} \approx -4.271400146 \end{cases}$$

[†] The asymptotic expression in (B.7) agrees with Zolotarev (1961).

APPENDIX C
Data of Coleman's Gats Experiment

H ₁		H ₂		V ₁		V ₂		N ₁		N ₂	
elevation	φ	elevation	φ	elevation	φ	elevation	φ	elevation	φ	elevation	φ
0	41	17	41	9	3	72	6	6	85	-51	13
-25	-13	-1	4	-30	-86	30	79	-4	33	-23	-73
54	-4	-19	-16	-19	72	8	6	12	-1	-18	-20
-58	47	-53	-57	14	-28	5	-27	31	-81	38	59
-18	-65	-28	-10	41	43	-11	24	-36	-76	-68	-73
27	54	-47	-2	52	77	23	-66	30	-48	68	30
-10	12	-18	-54	-33	2	-41	28	5	31	27	-59
-20	15	16	-24	-20	-5	-35	62	-24	-80	-66	-87
-8	26	-17	23	-5	16	-2	69	44	84	-65	-37
-60	84	-73	14	-27	3	-7	20	-13	-10	-57	-1
9	51	17	49	2	71	30	44	-2	55	24	57
21	-26	10	67	54	-75	-8	37	49	38	-35	-40
-40	-23	22	68	-33	65	-51	14	22	64	-24	-12
-5	-70	-36	-56	-26	70	-62	-18	39	-81	-28	-14
-5	-54	3	-43	28	-24	12	49	16	-44	9	10
		-4	-68	46	53	7	88			16	45

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